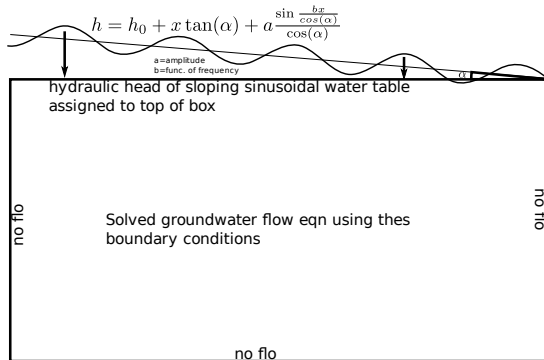


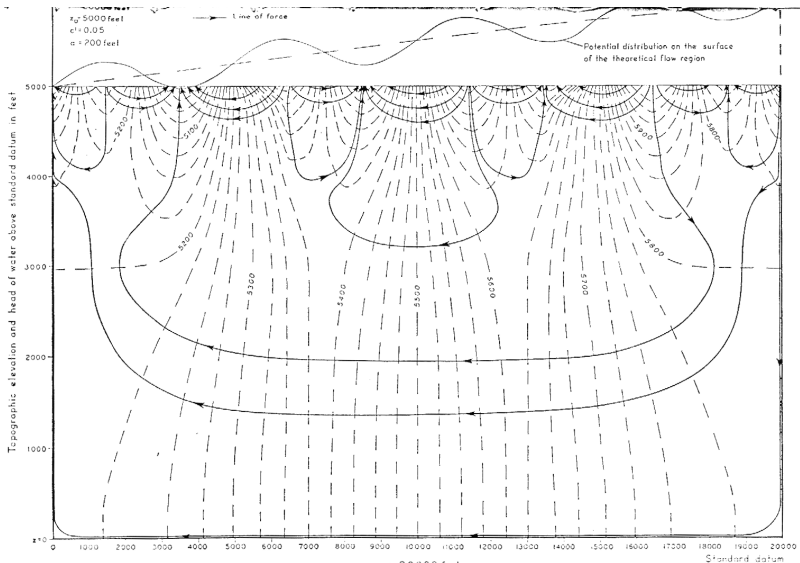
## Types of Groundwater Models

- ▶ Analog (eg. electrical,sandbox)
  - ▶ difficult to modify
  - ▶ difficult to make multilayer systems
- ▶ Analytic
- ▶ Numerical

# Analytic Models

- ▶ Classic Paper by Toth (1963)
- ▶ Solved groundwater flow equation for idealized basin





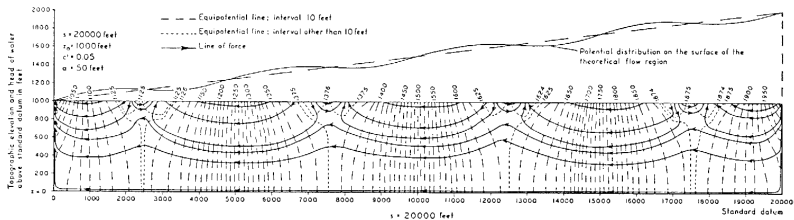


Fig. 2c. Potential distribution and flow pattern as obtained by equation 6.

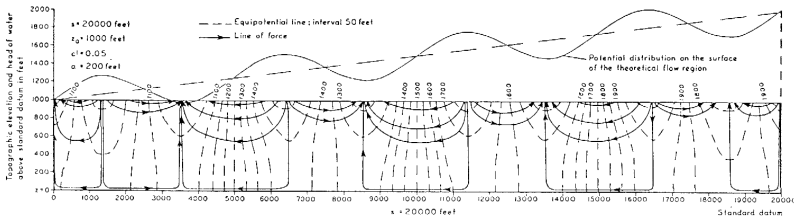
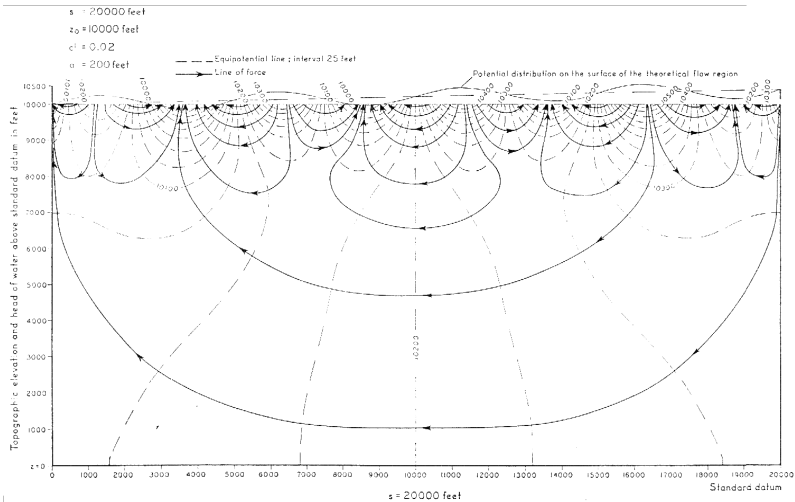


Fig. 2d. Potential distribution and flow pattern as obtained by equation 6.



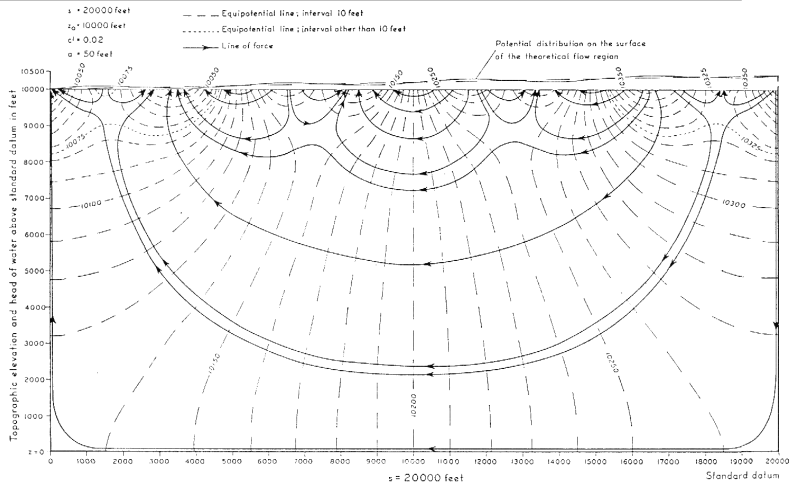


Fig. 2f. Potential distribution and flow pattern as obtained by equation 6.

# Current 'Types' of Computer Models

## Numerical Methods

- ▶ Finite Difference
- ▶ Finite Element
- ▶ Finite Volume
- ▶ Analytic Element

## Selecting a Model

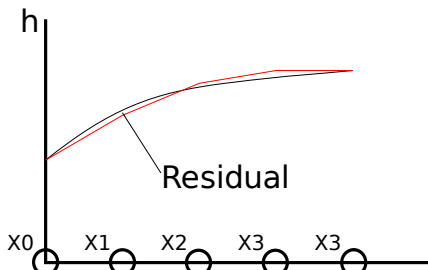
- ▶ 'Canned' Software
  - ▶ Modflow
  - ▶ Topodrive
  - ▶ Many others
- ▶ Program your own
  - ▶ Spreadsheet models
  - ▶ Software libraries
  - ▶ High level languages

# Finite Element Method

- ▶ Using weighting function (N) to interpolate across elements

$$\hat{h} = h_i \cdot N$$

$$\hat{h}_{01} = h_0 \frac{x - x_0}{\Delta x} + h_1 \left( 1 - \frac{x - x_0}{\Delta x} \right)$$





$$\frac{\partial^2 h}{\partial x^2} \rightarrow \frac{\partial^2 \hat{h}}{\partial x^2}$$

$$\frac{\partial^2 \hat{h}}{\partial x^2} = R$$

$$\int N \frac{\partial^2 \hat{h}}{\partial x^2} dx = \int N \cdot R dx = 0$$

integrate by parts

$$\sum N \frac{\partial h}{\partial x} - \sum h_i \int \frac{\partial N}{\partial x} \frac{\partial N}{\partial x} dx = 0$$

For each element:

$$\begin{bmatrix} \int_x \frac{\partial N_1}{\partial x} \frac{\partial N_1}{\partial x} & \int_x \frac{\partial N_1}{\partial x} \frac{\partial N_2}{\partial x} \\ \int_x \frac{\partial N_2}{\partial x} \frac{\partial N_1}{\partial x} & \int_x \frac{\partial N_2}{\partial x} \frac{\partial N_2}{\partial x} \end{bmatrix} \begin{bmatrix} h_1 \\ h_2 \end{bmatrix} = \begin{bmatrix} N_1 \frac{\partial h}{\partial x} \\ N_2 \frac{\partial h}{\partial x} \end{bmatrix}$$

note that  $\frac{\partial N}{\partial x} = \frac{1}{\Delta x}$  is a constant.

numerically integrate

$$\frac{1}{\Delta x_i} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \begin{bmatrix} h_1 \\ h_2 \end{bmatrix} = \begin{bmatrix} N_1 \frac{\partial h}{\partial x} \\ N_2 \frac{\partial h}{\partial x} \end{bmatrix}$$

left side applies to fluxes across boundary

Combining the terms from two elements:

$$\begin{bmatrix} \frac{1}{\Delta x} & -\frac{1}{\Delta x} & 0 \\ -\frac{1}{\Delta x} & \frac{1+1}{\Delta x} & -\frac{1}{\Delta x} \\ 0 & -\frac{1}{\Delta x} & \frac{1}{\Delta x} \end{bmatrix} \begin{bmatrix} h_1 \\ h_2 \\ h_3 \end{bmatrix} = \begin{bmatrix} \frac{\partial h}{\partial x_1} \\ 0 \\ \frac{\partial h}{\partial x_3} \end{bmatrix} \quad (1)$$

## Finite Difference Method

- ▶ truncated Taylor series used to approximate terms in PDE

$$h(x_0 + \Delta x) = \sum_{n=0}^{\infty} \frac{d(h(x_0))}{dx} \frac{\Delta x^n}{n!}$$

$$h(x_0 + \Delta x) = h(x_0) + \frac{dh(x_0)}{1!dx} (\Delta x) + \frac{d^2h(x_0)}{2!dx^2} (\Delta x)^2 \dots$$

$\frac{\partial h}{\partial t}$  can be approximated by truncating the Taylor series after the first term:

$$h(x + \Delta x) = h(x) + \frac{1}{1!} \left( \Delta x \frac{\partial h(x)}{\partial x} \right) + \left( \frac{\Delta x^2}{2} \frac{\partial^2 h(x)}{\partial x^2} \right) \dots$$
$$\frac{\partial h}{\partial x} = \frac{h(x + \Delta x) - h(x)}{\Delta x} + \left( \frac{\Delta x}{2} \frac{\partial^2 h(x)}{\partial x^2} \right) \dots$$

Truncating the after the second term results in an error proportional to  $\Delta x$  (first order). When  $\Delta x$  is small, the error will be small.

The forward Taylor series approximation is:

$$\frac{\partial h}{\partial x} \approx \frac{h(x + \Delta x) - h(x)}{\Delta x}$$

As similar procedure is followed to develop the backwards Taylor series approximation:

$$\frac{\partial h}{\partial x} \approx \frac{h(x) - h(x - \Delta x)}{\Delta x}$$

$$h(x + \Delta x) = h(x) + \frac{\Delta x}{1} \frac{\partial h(x)}{\partial x} + \frac{\Delta x^2}{2} \frac{\partial^2 h(x)}{\partial x^2} + \frac{\Delta x^3}{3!} \frac{\partial^3 h(x)}{\partial x^3} \dots$$

$$h(x - \Delta x) = h(x) - \frac{\Delta x}{1} \frac{\partial h(x)}{\partial x} + \frac{\Delta x^2}{2} \frac{\partial^2 h(x)}{\partial x^2} - \frac{\Delta x^3}{3!} \frac{\partial^3 h(x)}{\partial x^3} \dots$$

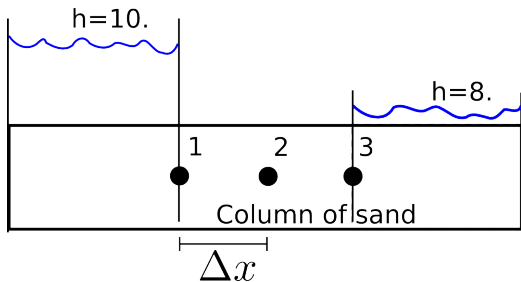
Adding these together yields:

$$h(x + \Delta x) + h(x - \Delta x) = 2h(x) + 2 \frac{\Delta x^2}{2} \frac{\partial^2 h(x)}{\partial x^2} + 2 \frac{\Delta x^4}{4!} \frac{\partial^4 h(x)}{\partial x^4} \dots$$

$$\frac{h(x + \Delta x) - 2h(x) + h(x - \Delta x)}{\Delta x^2} = \frac{\partial^2 h(x)}{\partial x^2} + 2 \frac{\Delta x^2}{4!} \frac{\partial^4 h(x)}{\partial x^4} \dots$$

Truncating the higher order terms gives second order accurate approximation for the second derivative. What happens if we subtract above equations?

## One-dimension problem I





## One-dimension problem II

For a 1-D, homogeneous, isotropic ground-water flow system:

$$0 = \frac{\partial^2 h}{\partial x^2}$$

Applying the finite-difference approximation to this problem results in an algebraic equation of each node in the problem area.

$$\begin{aligned}h_1 &= 10 \\ \frac{h_1 - 2h_2 + h_3}{\Delta x^2} &= 0 \\ h_3 &= 8\end{aligned}$$

What is  $h_2$ ? Note that the finite-difference method solves the problem at discrete points typically spaced at equal intervals.

# Selecting a modeling platform

- ▶ User needs
  - ▶ Ease of use
  - ▶ Versatility
  - ▶ Cost
- ▶ Software sources
  - ▶ USGS (MODFLOW, SUTRA and others)
  - ▶ USEPA (WHAEM, MODFLOW Support)
  - ▶ Commercial (FEFLOW, MIKE SHE, PORFLOW)
  - ▶ General CFD packages (OpenFoam, FiPy,...)

# Modflow

- ▶ Finite Volume (Integrated FD)
- ▶ Modular Structure
- ▶ Free from USGS
- ▶ Commercialized versions, GUI

## Streamlines in flownets

- ▶ Streamline position can be calculated using the stream function.
- ▶ In an isotropic system with an incompressible fluid,

$$\frac{dx}{q_x} = \frac{dy}{q_y} \quad (2)$$

- ▶ relates the length of a vector to the discharge.
- ▶  $\Psi$  is the stream function, a parameter that has a constant value along a streamline.
- ▶ If true, change in  $\Psi$  with distance along a streamline must be zero:

$$\frac{d\Psi}{ds} = 0 = \frac{\partial\Psi}{\partial x} \frac{dx}{ds} + \frac{\partial\Psi}{\partial y} \frac{dy}{ds} \quad (3)$$

Combining these last two equations:

$$\frac{\partial \Psi}{\partial x} = -\frac{\partial h}{\partial y} \quad (4)$$

$$\frac{\partial \Psi}{\partial y} = \frac{\partial h}{\partial x} \quad (5)$$